Induced Matter Theory & Heisenberg-like Uncertainty Relations

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### **Abstract**

We show that a differential variant of the Heisenberg uncertainty relations emerges naturally from induced matter theory, as a sum of line elements in both momentum & Minkowski spaces.

# **Introduction**

 Just after the introduction of Kaluza-Klein theory in the 1920's [1,2], Campbell [3] and later Magaard [4] have argued that it is possible to embed a n-dimensional theory into a n+1 dimensional manifold. Today, a few workers are using this basic notion to investigate a wide range of theoretical topics, from brane worlds[5-7] and cosmology [8-11], to relativistic quantum mechanics [12-14]. Induced matter theory posits that matter is in fact a direct manifestation of a non-compactified, fifth spatial dimension [15].

 Since matter is intrinsically quantum-mechanical (QM), we argue here that the origin of QM uncertainties in energy-momentum and space-time, must ultimately originate in the fifth dimension, and are thus geometric in nature. As Moffat has recently argued [16], there exist completely dualistic descriptions of classical particle motion, such that one may construct all the mathematical elements of general relativity, e.g., line elements, affine connexions, curvature tensors, space-time tensors, etc., in momentum space, in a fashion exactly analogous to a psuedo-Riemannian space.

 Here, we extend this program, and argue that the propagation of a complex scalar field in a 5-D space-time continuum, leads to both a 5-momentum conservation condition and in turn, to a geometrical origin of the Heisenberg uncertainty relations.

Index convention shall be Greek for  $3+1$  dimensions (e.g.,  $\mu = 0 - 3$ ) and Latin for 4 +1 dimensions (e.g.,  $A = 0 - 4$ ), with metric signatures (+---) and (+----) respectively.

#### Theory

We begin by investigating the propagation of a massive, complex scalar field, via the Klein-Gordon equation (KGE),

(1)  $\partial^{\mu} \partial_{\mu} \Psi + k_c^2 \Psi = 0$ , Where  $k_c = \frac{mc}{\hbar}$  is the Compton wavenumber & *m* is the mass of the scalar field quantum. In 5-D, the KGE goes over to the null D'Alembertian,

(2)  $\partial^A \partial_A \Psi = 0$  in which the Compton wavenumber is the eigenvalue of the 5<sup>th</sup> dimensional term  $\partial^l \partial_l \Psi = -k_{\rho}^2 \Psi$ . Representing the field as a plane wave, with  $l \partial_l \Psi = -k_c$  $\Psi = \exp(iS/\hbar)$ , and identifying S as the classical action, eq.(2) then becomes,

$$
(3) \quad \partial^A S \partial_A S + i\hbar \partial^A \partial_A S = 0
$$

Defining the 5-momenta as  $P_A = \partial_A S$ , we rewrite (3) as,

- Since the real & imaginary parts of this invariant must A  $P^A P_A + i\hbar \partial^A P$ vanish individually, eq.(4) must have a null real part,
- Expanding the sum gives the energy-momentum relation (5)  $P^A P_4 = 0$ of special relativity,
- (6)  $P^{\mu}P_{\mu} = P^{l}P_{l} = (mc)^{2}$  Similarly, the imaginary part can be expanded & written as

(7)  $\partial^{\mu}P_{\mu} = \partial^l P_l = \eta$  Where  $\eta = \frac{(mc)^2}{\hbar}$  is a constant required for dimensional consistency.

Rewriting eq. (7), and forming a bilinear null invariant,

- (8)  $(dP^{\mu} \eta dx^{\mu})(dP_{\mu} \eta dx_{\mu}) = 0$ , Expanding & solving for the cross-term, gives
- (9)  $dx^{\mu}dP_{\mu} = \frac{1}{2\eta} (dP^{\mu}dP_{\mu} + \eta^2 dx^{\mu}dx_{\mu})$  $\mu$  $\mu$  $\mu$  $\mu$  $^{\mu}dP_{\mu}=\frac{1}{2m}(dP^{\mu}dP_{\mu}+\eta)$ η  $dx^{\mu}dP_{\mu} = \frac{1}{2\pi} (dP^{\mu}dP_{\mu} + \eta^2 dx^{\mu}dx_{\mu})$ , redefining the two terms on the RHS as

dimension-less, differential invariants,

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(10) 
$$
d\rho^2 = \frac{dP^{\mu}dP_{\mu}}{(mc)^2}
$$
, and 
$$
d\chi^2 = k_c^2 dx^{\mu} dx_{\mu}
$$
 we can write eq.(9) as

(11) 
$$
dx^{\mu}dP_{\mu} = \frac{\hbar}{2}(d\rho^2 + d\chi^2)
$$

Eq.(11) is a Heisenberg uncertainty-like relation, in which the mean value of the two invariant line elements in canonically-conjugate metric spaces, contributes to a total uncertainty in four-momentum and -position, scaled by Dirac's action constant.

# Conclusion

From the physics of STM theory, we have shown that the propagation of a complex scalar field in 5-D, leads directly to an invariant with a form similar to the Heisenberg uncertainty relations. Moreover, the `backbone' of these relations originates in the line elements of the differential geometries of momentum and Minkowski spaces. This suggests that there may be a common descriptor to both the geometric and probability notions of relativistic quantum mechanics.

#### **References**

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